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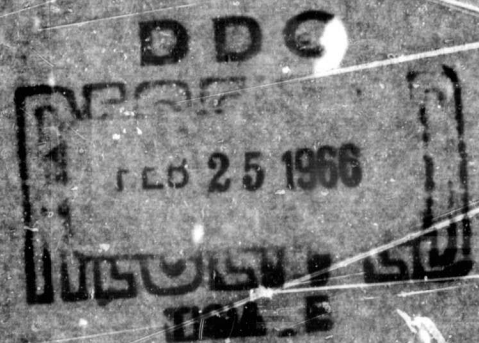
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EFFECTS OF MICROMETEOROID CRATERING  
ON THE DIRECTION OF THE AXIS  
OF MAXIMUM MOMENT OF INERTIA

H. O. Barthel

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## Introduction

The use of an artificial earth satellite as a gyroscope for measurement of a possible effect of general relativity, described by Schiff,\* is being studied. Mirrors on the surface of a fully passive spinning satellite will be used to reflect solar light for measuring a predicted change of six seconds of arc in the spin axis direction during a one year experiment. Although long term changes in the angle between the mirrors and the satellite spin axis can be corrected for, short term changes during the course of a measurement (several weeks) can cause significant errors in the data.

Because of the high degree of inertial symmetry necessary for reduction of gravity-gradient torque, meteoroids which do not significantly disturb the direction of angular momentum have a significant effect, after damping, on the direction of the body axis and mirrors relative to the spin axis. The effect of meteoroids on the principle axes of inertia, which ultimately control the body axis, is discussed herein.

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\*L. I. Schiff, Proc. Natl. Acad. Sci. U.S., 46, 871 (1960)



# Effects of Micrometeoroid Cratering on the Direction of the Axis of Maximum Moment of Inertia

H. O. Barthel

This examination of the wandering of the direction of the axis of maximum moment of inertia will proceed in the following way: First, the inertia tensor will be used to determine the effect of the loss of a small mass from the surface. The rotation of axes necessary to yield the principal moments of inertia will be found. Next, the cratering mechanism will be examined briefly and will lead to a consideration of the velocity distribution of micrometeoroids. Having obtained the velocity distribution, an estimate of the fraction of the total flux at a given mass having these velocities will be made. This information will then be used to estimate the number of hits per year, each of which could cause an angular shift of the order of one-tenth of the Schiff precession angle. Finally, a brief discussion of the results will be made.

The inertia tensor for an axis system  $x, y, z$  aligned in the direction of the principal moments of inertia of the rotationally symmetric body will be assumed to be

$$\tilde{I} = \begin{pmatrix} A & 0 & 0 \\ 0 & A & 0 \\ 0 & 0 & C \end{pmatrix} \quad (1)$$

before the loss of the mass  $m_t$  at the point  $x_s, z_s$ . After the loss of the mass  $m_t$ , it becomes

$$\tilde{I}' = \begin{pmatrix} A - m_t z_s^2 & -m_t x_s z_s & 0 \\ 0 & A - m_t (x_s^2 + z_s^2) & 0 \\ -m_t x_s z_s & 0 & C - m_t x_s^2 \end{pmatrix} \quad (2)$$

which can be made principal by the transformation

$$\tilde{B} = \begin{pmatrix} \cos \alpha & 0 & \sin \alpha \\ 0 & 1 & 0 \\ -\sin \alpha & 0 & \cos \alpha \end{pmatrix} \quad (3)$$

where  $\alpha$  is found from

$$\tan 2\alpha = \frac{2m_t x_s z_s}{C - A - m_t (x_s^2 - z_s^2)} \quad (4)$$

Equation 4, for an oblate spheroid for which  $C = \frac{2Ma^2}{5}$ ,  $A = \frac{M}{5}(a^2 + c^2)$

$c = a\sqrt{1 - e^2}$ , becomes

$$\tan 2\alpha = \frac{10 \frac{m_t}{M} \frac{x_s}{a} \sqrt{1 - \left(\frac{x_s}{a}\right)^2} \sqrt{1 - e^2}}{e^2 - 5 \frac{m_t}{M} \left\{ 2 \left(\frac{x_s}{a}\right)^2 - 1 + e^2 \left[ 1 - \left(\frac{x_s}{a}\right)^2 \right] \right\}} \quad (5)$$

For small  $\alpha$ ,  $\frac{m_t}{M} \ll e^2$ ,  $\frac{C-A}{C} = \frac{e^2}{2}$ , and equation 5 reduces to

$$\alpha = \frac{2.5 \frac{m_t}{M} \frac{x_s}{a} \sqrt{1 - \left(\frac{x_s}{a}\right)^2} \sqrt{1 - e^2}}{\frac{C-A}{C}} \quad (6)$$

Equation 6 will be assumed to hold. It can be solved for the mass ratio  $\frac{m_t}{M}$  to yield

$$\frac{m_t}{M} = \frac{\left(\frac{C-A}{C}\right) \alpha}{2.5 \left(\frac{x_s}{a}\right) \sqrt{1 - \left(\frac{x_s}{a}\right)^2} \sqrt{1 - e^2}} \quad (7)$$

which has a minimum value when  $\frac{x_s}{a} = .707$ . Substitution of this value of  $x_s / a$  and assuming  $\epsilon \ll 1$  results in

$$\frac{m_t}{M} \Big|_{\min} \doteq 0.8 \left( \frac{C-A}{C} \right) \alpha. \quad (8)$$

The mass  $m_t$  lost is dependent upon the hypervelocity impact of a micrometeoroid particle of mass  $m_p$ . The micrometeoroid has kinetic energy  $\frac{m_p V_p^2}{2}$ . This energy is divided between 1) the melting of the mass  $m_t$  in a hemispherical crater, 2) the kinetic energy of the mass  $m_t$  as it leaves, 3) radiation, 4) the raising of the temperature of the target around the crater, and 5) the generation of waves propagating in the target. In equation form, the melting of  $m_t$  can be stated as

$$m_t \epsilon = \eta_1 \frac{m_p V_p^2}{2} \quad (9)$$

where  $\epsilon$  is the energy necessary to melt a unit mass of the target and  $\eta_1$  is the fraction of the total kinetic energy of the impacting particle going into melting. Values of  $\eta_1$  are not known. Consequently, values will be assumed when calculations are made.

In order to use equation 9, values of  $V_p$  must be known. In much of the literature, it is assumed to be 30 km/sec. However, the distribution of velocities of meteoroids in Fig. 2 of NASA TN D1105 (Davison and Winslow, Space Debris Hazard Evaluation) has led to the following model. Of all particles colliding with the earth, approximately 40% are in retrograde orbits and 60% are in direct orbits. This difference is caused by a decrease in the earth's effective capture cross section as the relative speed increases.



Of those in retrograde orbits, approximately 45% are clustered near a velocity of 38 km/sec relative to the earth and 55% are clustered near 68 km/sec relative to the earth. These groupings arise from the likelihood that, over several billion years, the earth has swept up most of the particles having velocities relative to the sun of 30 km/sec, for these have circular orbits at the earth's distance from the sun. Of those in direct orbits, approximately one-third overtake the earth and have velocities near 15 km/sec relative to the earth and the other two-thirds are overtaken by the earth and have velocities near 22 km/sec relative to the earth.

Letting  $\Phi_i$  be the flux of particles having masses larger than  $m_p$  [calculated from  $\Phi_i = 10^{-17} (m_p)^{-1.7}$ , where  $m_p$  is in grams,  $\Phi_i$  is in number per square meter per second] and having velocity  $V_{p_i}$ , there results from the preceding considerations

$$\Phi = 0.2\Phi_{15} + 0.4\Phi_{22} + 0.18\Phi_{38} + 0.22\Phi_{68} . \quad (10)$$

The number of hits is then simply

$$N = \Phi A T S_f \quad (11)$$

where  $A$  is the surface area,  $T$  is the time and  $S_f$  is the shielding factor of the earth.

For a solid aluminum sphere ( $\epsilon = 250$  cal/gm) and assuming  $\alpha = .1\alpha_{\text{Schiff}} = 3.4 \times 10^{-6}$  rad,  $\eta_1 = 1$  and  $1/4$ ,  $T = 1$  year,  $S_f = .75$  (approximately 750 mile orbit) and  $A = 4\pi(.15)^2 = .283 \text{ m}^2$ , for  $\frac{C-A}{C} = 10^{-1}$ , and  $10^{-2}$  and  $10^{-3}$  there results

$\frac{C-A}{C}$	$N_{\eta_1 = 1}$	$N_{\eta_1 = 1/4}$
$10^{-1}$	$2.52 \times 10^{-2}$	$2.38 \times 10^{-3}$
$10^{-2}$	1.25	$1.18 \times 10^{-1}$
$10^{-3}$	62.8	5.91

These results are plotted in Fig. 1.

The assumption of  $\eta_1 = 1$  should yield an upper bound of hits per year for the assumed average flux levels. On the other hand, the value of  $\eta_1 = 1/4$  seemed to be a reasonable estimate in view of the mechanism of cratering. In any case, if a shower of 4 days duration with flux levels 100 times the average flux level occurs, then the yearly total will double.

The main contributions to the yearly total come from the retrograde micrometeoroids. If the distribution of velocities for these small particles is different, or the fraction of the total flux due to these particles is different, then the yearly totals will have to be revised.

In interpreting the results obtained above, one must keep in mind that each hit produces a maximum shift. For purely random hits, 75% of them will produce at least 1/2 of the maximum possible shift. Furthermore, for  $N$  greater than 25, there is a 50% chance of having a shift of 5 times the assumed shift. This result arises because in this number of hits one hit will involve a mass which is an order of magnitude larger than the assumed mass.

These considerations lead to the conclusion that the number of hits by masses which cause angular shift of .1 $\alpha$  Schiff should not

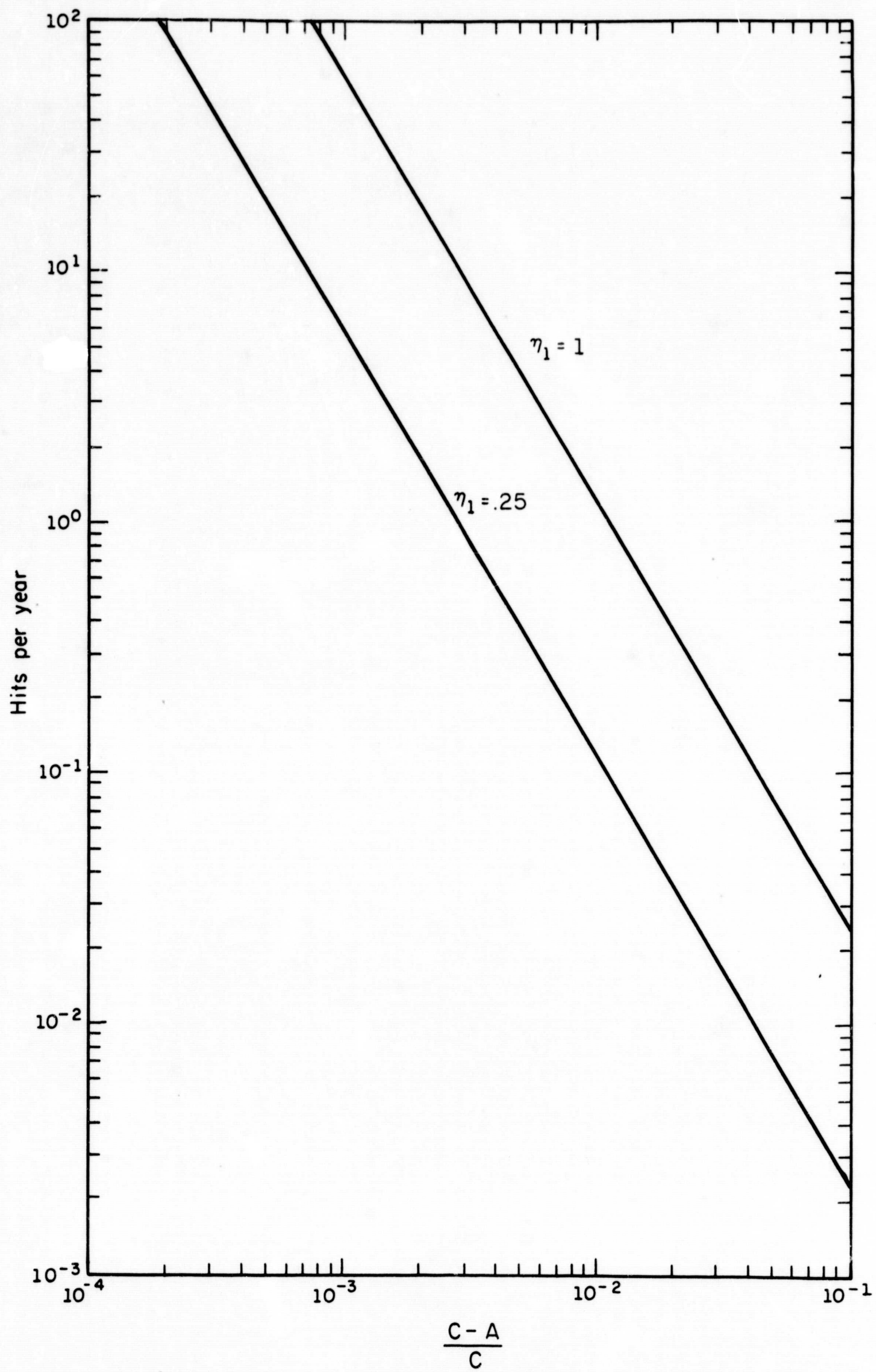


Fig. 1



exceed approximately 10 per year. This value also is likely to be reasonably consistent with the optical observation requirements. The detection of these angular shifts and the nutation decay should be accomplished before the next hit. These requirements indicate that  $\frac{C-A}{C}$  should be of the order of  $5 \times 10^{-3}$ .

If such a value of  $\frac{C-A}{C}$  cannot be used because of other considerations, then with a relatively large number of hits one is led to consider the shifts as a random walk problem. Let each angular shift from a hit be represented as a vector on a plane in which the origin represents the direction of the angular momentum vector and the location of the first point represents the direction of the axis of maximum moment of inertia at insertion. Then, the walk occurs in the following way. The first step will be from the first point toward the origin and its length will be dependent upon the time to the first hit and the nutation decay constant. The second step will be in a random direction and the length will depend upon the probabilities of 1) the striking mass being between  $m_p$  and  $m_p + dm_p$ , 2) the velocity of striking being between  $V_p$  and  $V_p + dV_p$ , and 3) the location and direction (relative to the normal to the surface) of the hit on the body is such that, in terms of the maximum shift for fixed  $m_p$  and  $V_p$ , the fractional shift will be between  $F$  and  $F + dF$ . The third step is a decay step toward the origin. The fourth step is similar to the second step. In the continuation of the walk, then all odd steps are toward the origin, all even steps are random. Consequently, one is led to probabilities of being certain distances from the origin after  $N$  hits. No attempt has been made to solve this problem.

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## KEY WORDS

## LINK A

## LINK B

## LINK C

ROLE

WT

ROLE

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ROLE

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 maximum moment of inertia  
 inertia tensor  
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